

The Silicon Vertex Detector of the Belle II Experiment

- ² J. Wiechczynski,^{r,*} K. Adamczyk,^r H. Aihara,^p S. Bacher,^r S. Bahinipati,^e
- ³ J. Baudot,^d P. K. Behera,^f S. Bettarini,^{j,k} T. Bilka,^b A. Bozek,^r F. Buchsteiner,^a
- G. Casarosa, *j*,*k* L. Corona, *k* S. B. Das, *g* G. Dujany, *d* C. Finck, *d* F. Forti, *j*,*k*
- ⁵ M. Friedl,^{*a*} A. Gabrielli,^{*l,m*} B. Gobbo,^{*m*} S. Halder,^{*i*} K. Hara,^{*q,n*} S. Hazra,^{*i*}
- 6 T. Higuchi,^o C. Irmler,^a A. Ishikawa,^{q,n} Y. Jin,^m M. Kaleta,^r A. B. Kaliyar,^a
- J. Kandra,^b K. H. Kang,^o P. Kodyš,^b T. Kohriki,^q R. Kumar,^h K. Lalwani,^g
- **K. Lautenbach**, ^c R. Leboucher, ^c J. Libby, ^f L. Martel, ^d L. Massaccesi, ^{j,k}
- ⁹ G. B. Mohanty,^{*i*} S. Mondal,^{*j*,*k*} K. R. Nakamura,^{*q*,*n*} Z. Natkaniec,^{*r*} Y. Onuki,^{*p*}
- ¹⁰ F. Otani,^o A. Paladino^A, ^{j,k} E. Paoloni, ^{j,k} K. K. Rao,ⁱ I. Ripp-Baudot,^d G. Rizzo, ^{j,k}
- Y. Sato,^q C. Schwanda,^a J. Serrano,^c T. Shimasaki,^o J. Suzuki,^q S. Tanaka,^{q,n}
- ¹² F. Tenchini, ^{*j*,*k*} R. Thalmeier, ^{*a*} R. Tiwary, ^{*i*} T. Tsuboyama, ^{*q*} Y. Uematsu, ^{*p*} L. Vitale, ^{*l*,*m*}

¹³ Z. Wang,^{*p*} H. Yin,^{*a*} L. Zani^{B,*c*} and F. Zeng^{*o*} (Belle-II SVD collaboration)

- ¹⁴ ^aInstitute of High Energy Physics, Austrian Academy of Sciences, 1050 Vienna, Austria
- ¹⁵ ^bFaculty of Mathematics and Physics, Charles University, 121 16 Prague, Czech Republic
- ¹⁶ ^cAix Marseille Université, CNRS/IN2P3, CPPM, 13288 Marseille, France, ^Bpresently at INFN Sezione di
- 17 Roma Tre, I-00185 Roma, Italy
- ¹⁸ ^dIPHC, UMR 7178, Université de Strasbourg, CNRS, 67037 Strasbourg, France
- ¹⁹ ^eIndian Institute of Technology Bhubaneswar, Bhubaneswar 752050, India
- ²⁰ ^f Indian Institute of Technology Madras, Chennai 600036, India
- ²¹ ^gMalaviya National Institute of Technology Jaipur, Jaipur 302017, India
- ²² ^hPunjab Agricultural University, Ludhiana 141004, India
- ²³ ^{*i*}Tata Institute of Fundamental Research, Mumbai 400005, India
- ²⁴ ^jDipartimento di Fisica, Università di Pisa, I-56127 Pisa, Italy, ^Apresently at INFN Sezione di Bologna,
- 25 I-40127 Bologna, Italy
- ²⁶ ^kINFN Sezione di Pisa, I-56127 Pisa, Italy
- ²⁷ ¹Dipartimento di Fisica, Università di Trieste, I-34127 Trieste, Italy
- ²⁸ ^mINFN Sezione di Trieste, I-34127 Trieste, Italy
- ²⁹ ⁿThe Graduate University for Advanced Studies (SOKENDAI), Hayama 240-0193, Japan
- ³⁰ ^oKavli Institute for the Physics and Mathematics of the Universe, University of Tokyo, Kashiwa 277-8583,
- 31 Japan
- ³² ^pDepartment of Physics, University of Tokyo, Tokyo 113-0033, Japan
- ³³ ^{*q*}High Energy Accelerator Research Organization (KEK), Tsukuba 305-0801, Japan
- ³⁴ ^rH. Niewodniczanski Institute of Nuclear Physics, Krakow 31-342, Poland
- 35 *E-mail:* wiechczynski@belle2.ifj.edu.pl

^{*}Speaker

[©] Copyright owned by the author(s) under the terms of the Creative Commons Attribution-NonCommercial-NoDerivatives 4.0 International License (CC BY-NC-ND 4.0).

36

The Belle II experiment operating at the asymmetric-energy e^+e^- SuperKEKB collider, located in Tsukuba (Japan), has been collecting data since March 2019. Its excellent vertexing abilities are provided by the vertex detector, part of which is the silicon-strip vertex detector (SVD) that plays a crucial role in the charged-particle tracking close to the interaction point. The SVD has operated successfully and efficiently over the whole period of data taking so far. In this article, we briefly discuss its purpose, structure and basic description of the front-end electronics. The main quantities related to the SVD performance are presented. The foreseen increase in SuperKEKB luminosity will lead to higher background, so we describe its impact on the SVD performance. A quick overview of the radiation damage campaign is presented to show the predicted behaviour of the sensors subjected to high radiation, whose level is constantly monitored. We also discuss the ongoing software development to account for the high occupancy expected in the future. In particular, the utilization of the SVD hit time information is presented as a very important quantity to suppress off-time background hits and tracks. Finally, the work done during the first long shutdown of SuperKEKB is briefly described, during which a major upgrade of the pixel detector has been successfully done. Resumption of the beam operation is expected in early 2024.

Keywords: Silicon strip detector, Vertex detector, Tracking detector, Belle II

The 32nd International Workshop on Vertex Detectors (VERTEX2023) 16-20 October 2023 Sestri Levante, Genova, Italy

37 1. Introduction

The Belle II [1] experiment is dedicated to search for physics beyond the standard model at 38 the intensity frontier. It operates at the SuperKEKB collider located at KEK, Tsukuba in Japan, 39 providing asymmetric beams of 7 GeV electrons and 4 GeV positrons. In the accelerator's default 40 operation regime, the center-of-mass energy is set to the $\Upsilon(4S)$ resonance, hence it produces a huge 41 sample of B mesons via the $e^+e^- \rightarrow \Upsilon(4S) \rightarrow B\bar{B}$ process. So far, SuperKEKB achieved the 42 highest instantaneous luminosity of 4.7×10^{34} cm⁻²s⁻¹, which is the current world record. The 43 Belle II detector is a multipurpose spectrometer characterized by excellent vertexing capability and 44 good hermeticity, which has accumulated 424 fb⁻¹ to date, and its final goal is to collect a data 45 sample of 50 ab^{-1} , that will be possible with a constant increase of the SuperKEKB instantaneous 46 luminosity up to our final goal of 6×10^{35} cm⁻²s⁻¹. 47

Belle II is composed of various sub-detectors with the vertex detector (VXD) being the closest 48 to the interaction point. It is divided into two further subsystems. The innermost part is the pixel 49 detector (PXD), which is based on depleted field effect transistor pixel sensors. The PXD consists 50 of two layers (numbered 1-2) and its main goal is the precise determination of the decay vertices. 51 Outside the PXD is the silicon-strip vertex detector (SVD) [2] with four layers (numbered 3-6) 52 that mostly extrapolates the measured tracks to the PXD, defining the so-called region of interest 53 (ROI), which significantly reduces the amount of data recorded by the PXD. The SVD also performs 54 standalone tracking for low-momentum charged particles and contributes to their identification by 55 providing energy loss information. 56

57 2. SVD structure

Each SVD layer is composed of a number of double-sided silicon strip detectors (DSSDs) 58 that are manufactured on an n-type bulk wafer with a thickness of about 300 μ m (Figure 1). One 59 side of the sensor is covered by the p-type silicon strips placed in parallel to the beam axis that 60 determine the $r - \phi$ coordinates (distance from the z-axis and azimuthal angle, respectively), and 61 the n-type strips are placed perpendicularly on the other side of the bulk, measuring the z coordinate 62 (collinear to the electron beam). Figure 1 (left) shows a schematic picture of SVD layers and 63 associated sensors with increasing numbering from the forward (FWD) to the backward (BWD) 64 regions. Such structure is repeated along the azimuthal direction forming different ladders and the 65 so-called windmill geometry of the SVD. The sensors differ depending on the layer and the region 66 in which they are placed in the SVD. In the FWD part, for layers 4-6, they have a trapezoidal shape 67

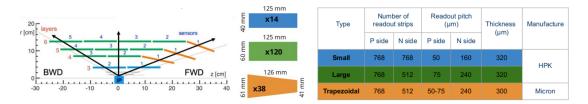


Figure 1: Schematic picture of SVD sensors forming different layers (left) and a table summarizing the parameters for each type of sensor (right).

and are slanted in the region that, due to the asymmetric beams, is characterised by the highest track multiplicity. In addition, in layer 3 the sensors are smaller and contain more n-type strips than that

- ⁶⁹ multiplicity. In addition, in layer 3 the sensors are smaller and contain more n-type strips than that ⁷⁰ in layers 4-6. This also implies the readout pitch (distance between two readout strips) to be much
- ⁷¹ smaller for p-side strips with respect to the n-side. To improve spatial resolution, a floating strip is
- ⁷² placed between two readout strips on both p- and n-sides. The charge induced in the floating strip
- ⁷³ is shared by the neighboring strips, reducing the effective strip pitch to half of the readout pitch.
- The right table of Figure 1 summarises the sensor parameters. The SVD consists of 224 thousand
- readout strips and 172 sensors with an active area of 1.2 m^2 .

76 2.1 Front-end electronics

For the readout we use APV25 chips [3]. For the central part of SVD (except for layer 3), the 77 chips are attached directly to the DSSD sensors via flex circuits bent over the DSSD edge (Origami 78 concept). The edge sensors use hybrid boards located outside the active volume. The APV25 has 79 128 channels per chip and amplifiers that provide a shaping time of 50 ns. Radiation hardness 80 exceeds 100 Mrad and the power consumption is around 0.4 W/chip. The sampling frequency is 32 81 MHz and after the trigger's arrival we can collect six consecutive signal samples in total with the 82 multipeak mode. To account for higher luminosity in the future, we have introduced the so-called 83 "3/6 mixed acquisition mode", which allows switching between three and six samples recorded on 84 an event-by-event basis, based on the trigger type (and hence its time accuracy) for a particular 85 event. This mode, already prepared and tested, significantly reduces the data size, which can be 86 crucial in high background conditions. 87

3. SVD performance

Since the start of the operation we have observed very smooth performance of the SVD, with a very few masked strips (less than 1%). Moreover, the environment has been stable and the evolution of calibration constants is consistent with expectation. Also, the effects of radiation damage are well under control.

Several quantities related to the SVD performance - sensor efficiency, signal-to-noise ratio, and 93 both spatial and time resolution - are constantly monitored. Regarding SVD sensor efficiency, the 94 values for all sensors are typically over 99% and they are also very stable over the whole period of 95 data taking. Clusters are formed from adjacent strips with significant signal and the charge collected 96 in a given cluster strongly depends on the incident angle of the track. Over time, we observe very 97 similar cluster charge in all the sensors once normalized to the track's length. For layer 4-5-6 on 98 the n-type strips we observe 10-30% loss of the signal due to the large pitch combined with the 99 presence of a floating strip. Another important quantity is the signal-to-noise ratio (SNR), which is 100 satisfactory for all 172 sensors. The SNR MPV is ranging from 13 to 30, depending on the sensor 101 position, due to the track incident angle with the sensor, and on the sensor side, with smaller SNR 102 for the p-sides, due to larger noise for the longer strip length. A small decrease of cluster SNR value 103 is observed in 2022 measurement, due to increased noise from radiation damage by approximately 104 20%– 30%. In Figure 2 the distributions of cluster charge (left) and SNR (right) are presented, 105 where histograms representing the data collected in 2020 and 2022 are superimposed. 106

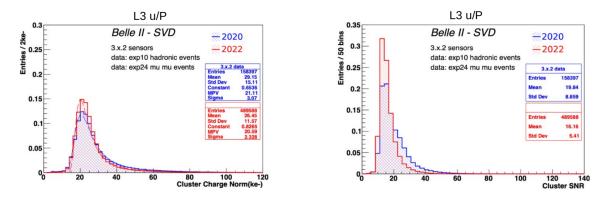


Figure 2: Distribution of cluster charge (left) and signal-to-noise ratio (right) for layer 3 (p-side). Comparison between data taken in 2020 (blue) and 2022 (red) is presented.

Both position and time resolution are very important metrics for excellent SVD performance. 107 The position resolution measurement is based on the residuals, i.e., the clusters' positions with 108 respect to the intercept of the unbiased tracks' extrapolation, and it is evaluated with a large sample 109 of $e^+e^- \rightarrow \mu^+\mu^-$ decays. As shown in Figure 3, this quantity depends on the incident angle and is 110 very stable during the period of the Belle II operation. Cluster position is calculated as the center 111 of gravity (i.e. weighted mean) of the various strip positions inside the cluster, using the collected 112 charge as strip weight. The cluster position resolution then depends on cluster size and on the 113 strip signal-to-noise ratio, varying with the track incident angle with an expected non monotone 114 dependence visible in Figure 3. For zero incident angle, with perpedicular tracks more likely to 115 produce a single strip cluster, one can only achieve a "digital" resolution, corresponding to the 116 standard deviation of a flat distribution inside implanted strip pitch (floating strip pitch/ $\sqrt{12}$). The 117 resolution improves with more than one strip in the cluster, adding the the center of gravity analog 118 information. It reaches a minimum for two strip clusters, corresponding to the incident angle for 119 which the projection of the track along the direction perpendicular to the strips on the detector 120 plane is two floating strip pitches (4 and 14 degrees for the layer 3 p- and n-sides pitches). For 121 large incident angle, since the signal collected in each strip decreases with the incident angle θ 122 (proportional to pitch/sin(θ)), the consequent reduction of the strip signal-to-noise ratio degrades 123 the cluster position resolution, visible in Figure 3. In general, the resolution for the n-side (left plot) 124 is about two times worse with respect to that for the u-side, which is a result of the different pitch 125 and strip signal-to-noise ratio on the two sides. 126

Hit time resolution is measured with respect to the event time of the collision provided by 127 central drift chamber (CDC) and exhibits a very good resolution of less than 3 ns for the clusters 128 associated to tracks. Using the average value of all the hits on a given track, the so-called "track-129 time" can be computed, slightly improving the time resolution. Furthermore, the "event-time" 130 can be determined using all the clusters associated to selected tracks in an event. In such a way, 131 the "event time" can be computed by the SVD with a resolution of the order of 1 ns, while the 132 computation is around 2000 times faster than the one based on CDC. This feature will be especially 133 important in the higher luminosity environment, as it can significantly speed up the reconstruction 134 process at the high-level trigger. 135

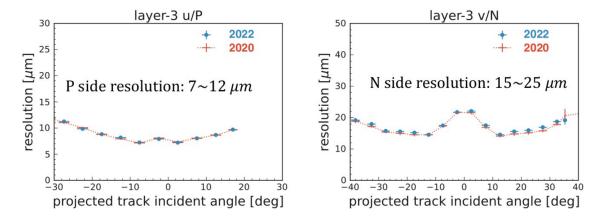


Figure 3: Distributions of position resolution for p-side (left) and n-side (right) as a function of the track incident angle. A comparison between data taken in 2020 (dots) and 2022 (dotted lines) is presented.

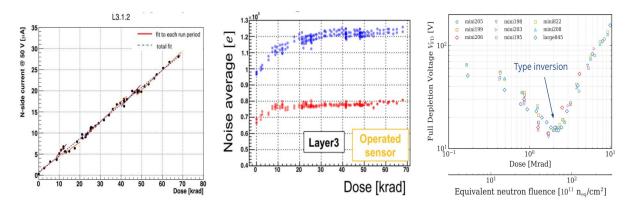


Figure 4: Left plot: Leakage current as a function of the accumulated dose; Center plot: the average noise level as a function of the accumulated dose for the p-side (blue dots) and n-side (red dots); Right plot: full depletion voltage as a function of the accumulated dose with the type inversion observed at 2 Mrad.

136 4. Radiation effects

In the high-energy physics experiments, the effects from radiation damage coming from ma-137 chine related background is a major factor that degrades the sensor performance with time. The 138 SVD accumulated dose is constantly measured using data from diamond sensors that are mounted 139 on the IP beam pipe, and the corresponding level of the equivalent neutron fluence is evaluated using 140 the ratio of equivalent neutron fluence to dose estimated from Monte Carlo simulation. Several 141 effects related to radiation damage must be taken into account. A linear increase of the leakage 142 current as a function of radiation damage is observed in the sensors, as expected from the bulk 143 damage described by the NIEL model [4], and shown in Figure 4 left. The sensor current is shown 144 as a function of the accumulated dose for one of the layer 3 sensors most exposed, that received 145 about 70 krad to date, corresponding to an equivalent neutron fluence of about $1.6 \times 10^{11} n_{ea}/cm^2$. 146 So far, this increase has had a negligible contribution to the noise because of both the small leakage 147

current and the short APV25 shaping time. The rate of the leakage current increase measured 148 is consistent with the experience from other experiments working with similar detectors and in 149 comparable conditions [5]. However, we expect some significant impact on the strip noise due to 150 the sensor leakage current, and hence a deterioration in SNR, for the dose of ~ 6 Mrad, which is 151 considered as SVD dose limit to preserve optimal performance. The strip noise for unirradiated 152 modules is dominated by the interstrip capacitance. During the operation we have observed an 153 increase in its value of about 20% (30%) for n-side (p-side), due to effects of surface radiation 154 damage that increases the interstrip capacitance, but it is expected to saturate, as also visible in 155 Figure 4 center. 156

Another relevant effect of the bulk radiation damage is the impact on depletion voltage. 157 The expected future radiation levels at the nominal luminosity, of about 0.35 Mrad/year and $8 \times$ 158 10^{11} n_{eq}/cm²/year, are affected by large uncertainty due to the machine evolution as well as a 159 possible redesign of the interaction region. To better explore the possible effects of bulk damage in 160 the SVD sensors after bulk type inversion, an irradiation campaign was conducted in July 2022 at 161 ELPH, Tohoku University. Several SVD sensors have been exposed to a 90 MeV electrons beam, 162 up to 10 Mrad, corresponding to an equivalent neutron fluence of 3×10^{13} n_{ea}/cm². The decrease 163 of the depletion voltage has been observed up to the point of bulk type inversion, which occurred at 164 2 Mrad (~ $6 \times 10^{12} n_{eq}/cm^2$), after which the depletion voltage started to increase again (Figure 4 165 right). Detailed measurements, whose results will be shortly published, confirmed that the sensors 166 will still work fine after the type inversion, which meets our expectation for these types of silicon 167 detectors. Since the beginning of the detector operation, we have not observed any change in the 168 depletion voltage in the sensors installed in the SVD, as expected due to the small accumulated 169 equivalent neutron fluence so far, below $2 \times 10^{11} n_{eq}/cm^2$. Considering all these results, the dose 170 limit of 6 Mrad and the extrapolation of the background levels quoted above, the SVD has a wide 171 safety margin for the accumulated radiation damage even after 10 years of the operation at the target 172 luminosity. 173

174 5. High background scenario and related software/hardware developments

With the increase of the luminosity and the expected larger machine related background, the 175 SVD occupancy will also increase and a deterioration of the tracking performance is expected above 176 certain levels. So far, the average hit occupancy is 0.5% for layer 3, which does not degrade the 177 performance. Nonetheless, the background extrapolation for different future scenarios has been 178 performed with detailed simulations of the various contributions (beam-gas, Toushek, etc.) and 179 applying appropriate data-simulation scale factors [6]. These studies predict that for the nominal 180 luminosity we can reach an occupancy in layer 3 very close to the limit of 4.7%, above which the 181 tracking performance deteriorates. These predictions have large uncertainties coming from poorly 182 known machine evolution in the future, with a possible redesign of the interaction region. In the 183 most conservative scenario, the layer-3 occupancy can increase up to $\sim 8.7\%$, which is far beyond 184 the modest tracking performance. Such a scenario motivates us to develop the SVD reconstruction 185 software, as well as to seriously consider the VXD upgrade [7], since the safety factor might be too 186 small to ensure good quality data. The technology assessment related to this hardware upgrade is 187 currently ongoing. 188

An important effort related to the software development is the utilization of the hit time 189 information from the SVD. The real signal hits come from well-triggered collisions, but the SVD 190 acquisition window (~ 100 ns) is much wider with respect to the SuperKEKB bunch spacing (6 ns). 191 Therefore, we need to cope with many off-time hits related to the beam-induced background or 192 background from the other bunches. The current selection is based on two requirements: a) time 193 difference between p- and n-side cluster, $|t_p - t_n| < 20$ ns, and b) the absolute value of the cluster 194 time, $|t_{p,n}| < 50$ ns. These criteria reject the majority of the background hits retaining above 99% of 195 the signal, and based on them the SVD occupancy limit for layer 3 can be set at 4.7%. Recently, a 196 more effective background suppression method has been developed in the form of so-called "SVD 197 grouping". It is based on an event-by-event classification of the clusters by their time, so the 198 clusters belonging to tracks from the same collisions are collected in the same group. Clusters 199 from the different collisions or beam background will be placed in the other groups; finally, only 200 the clusters belonging to the priority group will be used for the tracking. This feature reduces the 201 fake rate (fraction of the fake tracks) by 16% for the high-background scenario. An additional fake 202 rate reduction can be achieved by applying the selection on the track-time to reject off-time tracks. 203 Finally, these improvements allow an increase of the SVD occupancy limit for layer 3 from 4.7% to 204 around 6%. 205

206 6. Activities during the Long Shutdown 1

Long shutdown 1 (LS1) started in May 2022 and one of the goals was to upgrade the VXD with a 207 new PXD. During the first data taking period, the second layer of PXD was only partially equipped, 208 and 5/6 of the azimuthal angle remained uncovered. The new PXD provides the full coverage, 209 which is beneficial for more precise vertexing. Hardware activities for the VXD uninstallation and 210 reinstallation were intense: after the VXD extraction from Belle II, the SVD was detached from 211 the old PXD (May 16-17, 2023), then the new PXD was attached to the SVD (June 20-21, 2023) 212 and finally the complete VXD was installed in the Belle II detector. The whole delicate procedure 213 had neither major problems nor caused any damage. In the period of September 12 - October 1, 214 2023, the VXD commissioning was performed to confirm the PXD and SVD performance, and 215 also to check the impact from the increased PXD power consumption and possible increase in 216 the temperature on the sensor leakage current. From September 21, several cosmic runs with no 217 magnetic field were taken to check the performance and compare them with corresponding ones 218 for 2022 data samples. We observed no issues, in particular the noise distributions over readout 219 channels remained basically unchanged as well as SNR for the clusters associated to the tracks, with 220 stable excellent efficiency for all the sensors. 221

222 7. Conclusions

To conclude, SVD has successfully operated since March 2019 with very smooth performance and without major problems. Its good vertexing quality has been confirmed by many physics measurements, in particular those related to the lifetime analyses e.g. Ref. [8]. Some radiation damage effects were observed, but without any impact on the performance so far. However, the extrapolated background level indicates that the occupancy in the SVD can exceed
the current limit that guarantees good tracking performance. Hence, several software improvements
are being implemented to account for high background conditions. In particular, exploitation of
the SVD hit time is of major importance. Alongside, a VXD upgrade is also under discussion to
increase robustness against high background and to match a possible new interaction region.
The VXD reinstallation at Belle II with complete PXD has been successfully done during the
LS1, followed by successful VXD commissioning with cosmic data. The beam operation is planned

to resume in early 2024.

235 References

- [1] T. Abe et al., Belle II Technical Design Report, arXiv:1011.0352 (2010).
- ²³⁷ [2] K. Adamczyk et al., JINST **17**, P11042 (2022).
- [3] M. J. French et al., Nucl. Instrum. Meth. A 466, 359 (2001).
- ²³⁹ [4] G. Lindstrom et al., Nucl. Instrum. Meth. A **465**, 60-69 (2000).
- ²⁴⁰ [5] B. Aubert et al., Nucl. Instrum. Meth. A **729**, 615 (2013).
- [6] A. Natochii et al., Nucl. Instrum. Meth. A **1055**, 168550 (2023).
- ²⁴² [7] M. Babeluk et. al., Nucl. Instrum. Meth. A **1048**, 168015 (2023).
- [8] F. Abudinén et al., Phys. Rev. Lett. **130**, 071802 (2023)